Modelled trends in Antarctic sea ice thickness

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Abstract

Unlike the rapid sea-ice losses reported in the Arctic, satellite observations show an overall 1 increase in Antarctic sea ice concentration over recent decades. However, observations of 2 3 decadal trends in Antarctic ice thickness, and hence ice volume, do not currently exist. In this study a model of the Southern Ocean and its sea ice, forced by atmospheric reanalyses, is 4 used to assess 1992-2010 trends in ice thickness and volume. The model successfully 5 reproduces observations of mean ice concentration, thickness, and drift, and decadal trends in 6 ice concentration and drift, imparting some confidence in the hindcasted trends in ice 7 8 thickness. The model suggests that overall Antarctic sea ice volume has increased by approximately $30 \text{km}^3/\text{y}$ (0.4%/y) as an equal result of areal expansion ($20 \times 10^3 \text{km}^2/\text{y}$, or 9 0.2%/y) and thickening (1.5mm/y, or 0.2%/y). This ice volume increase is an order of 10 magnitude smaller than the Arctic decrease, and about half the size of the increased 11 freshwater supply from the Antarctic Ice Sheet. Similarly to the observed ice concentration 12 trends, the small overall increase in modelled ice volume is actually the residual of much 13 14 larger opposing regional trends. Thickness changes near the ice edge follow observed concentration changes, with increasing concentration corresponding to increased thickness. 15 Ice thickness increases are also found in the inner pack in the Amundsen and Weddell seas, 16 where the model suggests that observed ice-drift trends directed towards the coast have 17 caused dynamical thickening in autumn and winter. Modelled changes are predominantly 18 19 dynamic in origin in the Pacific sector and thermodynamic elsewhere.

21 **1. Introduction**

Arctic sea ice extent has declined rapidly in recent decades (-52×10^3 km²/y for 1979-2010), 22 but Antarctic sea ice extent has slowly increased $(+17 \times 10^3 \text{ km}^2/\text{y})$ over the same period 23 24 (Cavalieri and Parkinson 2012; Comiso and Nishio 2008; Parkinson and Cavalieri 2012; Zwally et al. 2002), raising fundamental questions of why the two poles have evolved so 25 differently in the context of climate change. The small overall Antarctic increase in ice area 26 is actually the residual of a coherent pattern of much larger regional increases and decreases 27 that almost compensate each other. These large local areal changes (up to 2% per year 28 29 increase and decrease, or 60% total over 30 years; Turner et al. (2009)) can also be viewed as changes in the length of the ice season (up to 3 days per year, or 3 months in total; 30 Stammerjohn et al. (2012)). The local changes are of the same magnitude as those in the 31 32 Arctic, which does not feature the regions of ice expansion that, in the Antarctic, more than offset the regions of loss. 33

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It is currently unclear exactly what causes the regional pattern of changes that produces the 35 overall increase in ice cover. Proposed drivers include changes in atmospheric temperature 36 or wind stress (Lefebvre and Goosse 2005; Liu et al. 2004; Turner et al. 2009), precipitation 37 (Liu and Curry 2010), ocean temperature (Jacobs and Comiso 1997), atmosphere and ocean 38 feedbacks (Stammerjohn et al. 2008; Zhang 2007), and increased freshwater flux from the 39 40 Antarctic Ice Sheet (Bintanja et al. 2013). Recent work has shown that the trends in Antarctic ice concentration are associated with trends in ice drift, and that both are caused by changes 41 in near-surface winds through a combination of dynamic and thermodynamic effects 42 (Holland and Kwok 2012). However, the ultimate cause of the relevant wind changes 43 remains uncertain. 44

The current generation of coupled climate models are unable to capture the increase in overall 46 Antarctic sea ice extent, instead hindcasting a decline in ice cover of a similar magnitude to 47 their modelled Arctic (Turner et al. 2012). This suggests that important deficiencies exist in 48 49 our understanding of ice and climate physics that will be relevant to the prediction of climate at both poles. The model projections of most aspects of Antarctic climate are questionable if 50 they cannot reproduce past observations of sea ice extent, since it is one of the better-51 monitored polar climate variables. Improved climate models are also required to answer top-52 level questions about past changes in Antarctic sea ice that are of vital importance to 53 54 policymakers. For example, it is unclear why Antarctic sea ice is not rapidly declining in response to increased greenhouse-gas concentrations and depletion of stratospheric ozone, 55 both of which are found to decrease Antarctic sea ice in coupled climate models (Bitz et al. 56 57 2006; Sigmond and Fyfe 2010). Several studies have suggested that the observed increase is an unlikely result of natural variability, which would consequently only be captured by a 58 small proportion of simulations (Mahlstein et al. 2013; Polvani and Smith 2013; Swart and 59 Fyfe 2013). However, this result is marginal, valid only for the annual-mean circumpolar 60 trend (Swart and Fyfe 2013), and relies upon the models having realistic natural variability, 61 which is not the case (Turner et al. 2012; Zunz et al. 2013). In this study we investigate the 62 Antarctic ice trends further in a quest to provide additional insight into these model 63 weaknesses. 64

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A critical gap in our understanding of Antarctic sea ice and its trends is caused by the relative paucity of Antarctic ice thickness data. Though spatially widespread, in-situ observations are severely lacking in spatial and temporal detail (Worby et al. 2008). Ice thickness can be determined from satellite altimetry by measuring the ice freeboard and assuming that the ice is freely floating with some choice of ice, snow, and seawater properties. Radar altimeters

have provided a relatively long record of ice thickness in the Arctic, but are subject to 71 variable snow penetration in the Antarctic that currently precludes the reliable determination 72 of ice freeboard (Giles et al. 2008). The freeboard of Antarctic ice and snow can be 73 accurately measured using laser altimeters, and this can be converted to ice thickness using 74 independent estimates of snow thickness and snow and ice density (Markus et al. 2011; Xie et 75 al. 2013; Zwally et al. 2008). Recent studies demonstrate that ice thickness can be derived to 76 a reasonable level of accuracy using the simple assumption that the snow—ice interface is at 77 sea level (Kurtz and Markus 2012; Ozsoy-Cicek et al. 2013). However, the available laser 78 79 altimeter data are limited in temporal coverage, and therefore unable to provide reliable trends in ice thickness. Instead, we use a coupled ice-ocean model to investigate the ice 80 thickness trends and their drivers. 81

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Models have previously been used to study various sensitivities of Antarctic sea ice, 83 including the effects of surface precipitation (Powell et al. 2005), winds (Stossel et al. 2011), 84 and ice-shelf meltwater (Hellmer 2004). Models have also been used to assess linkages 85 between sea-ice variability and large-scale climate modes (Lefebvre and Goosse 2005, 2008). 86 Several such models have been validated against ice observations, including those of ice 87 thickness, with notable success (Fichefet et al. 2003a; Losch et al. 2010; Timmermann et al. 88 2002; Timmermann et al. 2004; Timmermann et al. 2005; Timmermann et al. 2009). 89 However, only a few model studies consider changes in Antarctic sea ice thickness or 90 volume. 91

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Proposing an ocean feedback on increasing Antarctic sea ice, Zhang (2007) simulated a 1979-2004 increase in ice volume of 200 km³/y. The mean ice area of 10^7 km² thus implies a Antarctic-mean thickening of 2 cm/y, or 0.5 m over the period, which seems unfeasibly large.

This simulation had many shortcomings: overestimation of annual-mean ice volume by a 96 factor of 2 (Kurtz and Markus 2012); overestimation of area trend by a factor of 3; 97 disagreement with observed spatial pattern of concentration trends; disagreement with 98 observed temporal variability in total ice extent. Fichefet et al. (2003a) found an area 99 increase of 11×10^3 km²/y over 1958-1999 but no appreciable trend in ice thickness, though 100 considerable wind-driven decadal variability in ice thickness and area were identified. 101 Fichefet et al. (2003b) investigated 1955-2001 area and volume trends, finding an overall 102 decrease of 9×10^3 km²/y and increase of 11 km³/y respectively. However, reanalysis-forced 103 104 models should be treated with extreme caution prior to the onset of satellite sounding data assimilation in late 1978 (Bromwich and Fogt 2004), and the latter model produces no trend 105 in ice area during 1978-2001. Timmermann et al. (2005) report little modelled trend in ice 106 area or volume during 1977-1999, attributing this to their spin-up technique of repeating the 107 reanalysis forcing twice. Timmermann et al. (2009) model an ice area increase of 11×10^3 108 km^2/y for 1979-2006 but do not report the corresponding volume trend. Crucially, the latter 109 four studies do not show the spatial distribution of ice concentration trends, so it is impossible 110 to assess whether the physical processes driving their overall trends are realistic. 111

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A recent modelling study by Zhang (in press) specifically investigates the effect of changes in 113 winds on Antarctic ice volume. The study finds an increase in ice volume of 69 km^3/v over 114 1979-2010, but at 15×10^3 km³ the annual mean ice volume in this model is approximately 115 twice that inferred from observation (Kurtz and Markus 2012), which casts significant doubt 116 on the value of the volume trend. Since the ice extent is reasonable, this implies that the ice 117 thickness is approximately twice the true value. The study does not examine in detail the 118 changes during different seasons, in different regions, or the thermodynamic and dynamic 119 mechanisms underlying the changes. 120

Probably the most reliable estimate of recent ice volume trends are from the model of 122 Massonnet et al. (2013), which formally optimises the estimate by assimilating ice 123 concentration data using an ensemble Kalman filter. The results show an overall 1980-2008 124 increase in ice volume of 36 ± 34 km³/y, with a regional pattern of ice-thickness trends that are 125 closely related to the changes observed in ice concentration. The use of data assimilation has 126 strengths and weaknesses; the results should be quantitatively as reliable as possible, but the 127 adjustments made to the model state vector do not have a directly physical origin, and none 128 129 of the ice or ocean variables are conserved (Massonnet et al. 2013; Mathiot et al. 2012). This implies that the physical processes underlying any ice thickness changes cannot be examined. 130 Also, the need to run an ensemble of models limits the resolution possible in each case; 131 Massonnet et al. (2013) run 25 ensemble members at 2° resolution. 132

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The goal of this paper is to produce a high-resolution, free-running, observationally validated 134 hindcast of trends in Antarctic sea ice thickness and volume. This study is complementary to 135 those of Massonnet et al. (2013) and Zhang (in press); the results will not be quantitatively 136 perfect but the use of a free-running (non-data-assimilating) model ensures that thickness 137 trends are the result of calibrated model physics, which we examine in temporal and spatial 138 detail. We place particular emphasis on a detailed assessment of our model results against 139 satellite observations of the mean fields of Antarctic ice concentration, drift, and thickness, 140 and the trend fields of ice concentration and drift. This validation provides a clear view of 141 the relative confidence in the hindcasted ice thickness trends in different regions and seasons. 142

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144 **2. Methods**

We use revision c62r of the MITgcm (http://mitgcm.org) in a regional model of all ocean, sea 145 ice, and ice shelves south of 30°S. The ocean component solves the Boussinesq Navier-146 Stokes equations on a generalised curvilinear grid using an Arakawa C-grid finite-volume 147 discretisation and z-levels in the vertical (Marshall et al. 1997). All components use the same 148 mesh, with a locally isotropic horizontal resolution of 0.25° in longitude, producing 149 approximately square cells ranging from ~10 km on each side at 70°S to ~18 km at 50°S. 150 The ocean component has 50 vertical levels ranging from 10-m resolution over the top 100 m 151 to 457 m in the layer beneath 5461 m, though the step-like representation of seabed and ice-152 shelf topography is alleviated by the use of partial cells (Adcroft et al. 1997). Horizontal 153 diffusivity is parameterised following Gent and McWilliams (1990) with a variable 154 diffusivity (Visbeck et al. 1996) (limited to maximum 300 $\text{m}^2 \text{ s}^{-1}$) and slope-clipping (Large 155 et al. 1997). Horizontal viscosity is flow-dependent (Leith 1996). Vertical mixing is 156 parameterised according to the 'K-profile parameterisation' (KPP) scheme (Large et al. 157 1994), which combines representations of ocean internal mixing and the surface mixed layer, 158 exerting a significant influence upon the sea ice. A fully non-linear equation of state is used 159 (McDougall et al. 2003). 160

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The sea-ice component (Losch et al. 2010) is also formulated on a C grid. In this study we 162 use an Elastic-Viscous-Plastic procedure to solve for ice dynamics with an elliptic yield 163 164 curve. Free-slip conditions are applied at boundaries, and ice stress is applied directly to the surface of the ocean. Ice thermodynamics are treated using the 'zero layer' approach, 165 employing a constant thermal conductivity and linear temperature profile within the ice 166 (Semtner 1976). The model has only two prognostic ice classes (ice and water) but a linear 167 distribution of 7 thickness classes is used in the thermodynamic calculations. A prognostic 168 snow layer floods into ice if depressed below sea level. Ice salinity is neglected entirely, 169

which implies a slight over-prediction of freshwater fluxes because sea ice is in reality slightly saline. All prognostic variables are transported using first-order upwind advection. Far more sophisticated physical treatments of ice processes are available (Hunke and Lipscomb 2010), and it would be instructive to examine the effect of those in a future study, but the sea ice model is demonstrably able to reproduce the relevant ice observations (see below), so we are confident that its features are sufficient to support the conclusions of this study.

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178 Initial conditions for ocean temperature and salinity are taken from the World Ocean Atlas (Boyer et al. 2009) (extrapolating southwards where required) and seabed and ice-shelf 179 topography is taken from the RTOPO dataset (Timmermann et al. 2010). Steady 180 181 climatological boundary conditions are applied at 30°S, with temperature and salinity taken from the World Ocean Atlas and ocean velocities taken from the ECCO2 reanalysis 182 (Menemenlis et al. 2005). The ocean and sea-ice surfaces are forced using 6-hourly fields 183 from the ERA-Interim reanalysis (Dee et al. 2011) at a resolution of 1.5° in both longitude 184 and latitude. The forcing variables consist of zonal and meridional 10-m winds, 2-m air 185 temperature and specific humidity, downward shortwave and longwave radiation, air pressure 186 loading and precipitation. The pressure loading and thermodynamic interactions of static ice 187 shelves are also included (Losch 2008). Iceberg melting is a significant source of freshwater 188 189 to the Southern Ocean that occurs in a heterogeneous pattern depending upon the distribution of the bergs. We experimented with deriving this flux from third-party model fields, but 190 these were completely dependent upon the modelled bergs and could never be truly 191 representative of the time period used. Therefore, iceberg melting was represented simply by 192 distributing a freshwater flux of 2000 Gt/y uniformly around the coast (Jacobs et al. 1992). 193 No ocean salinity restoring is used. 194

The paucity of in-situ atmospheric data over the Southern Ocean means that reanalysis 196 forcing data contain significant biases prior to the onset of satellite sounding data assimilation 197 in late 1978 (Bromwich and Fogt 2004). Therefore, the model is first spun-up by repeating 198 1980 forcings 10 times, and then run forward from 1981 until the end of 2011. Starting the 199 simulations in January avoids the need for any initial sea-ice distribution. Validation of the 200 model against observed ice trends is essential to impart confidence in the modelled ice 201 thickness trends, so we analyse only the period 1992–2010, for which reliable data of trends 202 203 in ice concentration and drift are available (Holland and Kwok 2012). This provides a total of 22 years of model spin-up time, and we are confident that the trends presented are the 204 result of the atmospheric forcing, not ocean adjustment from initial conditions. In particular, 205 206 a test simulation in which the 1980 forcings were repeated for 40 years shows no significant sea-ice trends after year 20. 207

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The model validation requires observations of ice variables on an Antarctic-wide scale. Such 209 observations do not exist directly, but can be derived from quantities observable by satellite. 210 Daily ice concentration data generated from passive microwave emissions using the 211 Bootstrap algorithm are used, with all values below 0.15 masked (Comiso 2000). Ice drift 212 data generated by feature-tracking in the same passive microwave data are also available 213 214 daily for the entire period, though only from April—October due to a high rejection rate of data in the Austral summer (Holland and Kwok 2012; Kwok et al. 1998). The only 215 comprehensive ice thickness data available on the Antarctic-wide scale are from ICESat laser 216 altimetry campaigns, covering 1—3 one-month-long periods per year for 2003—2008 (Kurtz 217 and Markus 2012). These ice- and snow-thickness data are derived from measurements of 218 freeboard and the assumption that the snow-ice interface is at sea level; i.e. that all 219

freeboard is snow and all draft is solid ice. This assumption is highly questionable in detail,
but appears to provide a reasonable level of agreement with in-situ observations overall
(Kurtz and Markus 2012; Ozsoy-Cicek et al. 2013; Worby et al. 2008).

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Various definitions of 'ice thickness' are used in the literature. Throughout this study, 224 'effective ice thickness' is defined as the volume of ice per unit area of ocean, which is the 225 quantity conserved by the model, while 'average ice thickness' is used to refer to the volume 226 of ice per unit area of ice, which is closer to the quantity measured in the field. Effective ice 227 228 thickness is the product of the average ice thickness and the ice area concentration. We generally investigate fields of effective ice thickness because that is the quantity most 229 relevant to the overall changes in ice volume, but the Antarctic-wide average ice thickness is 230 231 also examined. We consider the thickness of ice only, rather than including the ice-borne snow layer, because the ice component is of greater interest to many scientific questions, and 232 is also better constrained in our model, which uses uncertain reanalysis precipitation fields to 233 generate ice-borne snow. We consider seasonal maps of means and trends calculated from 234 monthly-mean model output. Mean fields for each season are the overall average of all 235 appropriate months from all years. To produce trend fields, for each grid point we first 236 convert the model output into a timeseries of season-mean values, and then calculate the 237 interannual trend for each season from the appropriate seasonal values over the different 238 239 years. For example, to calculate the trend in winter ice concentration, we create fields of the mean ice concentration for each winter and then plot, at each grid point, the interannual 240 trends in those fields. 241

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243 **3. Results**

Before examining our results it is worth considering the extent to which we would expect real 244 ice thickness trends to be represented in a free-running hindcast model. In any model forced 245 by atmospheric reanalyses, ice extent (the ocean area covered by an ice concentration of at 246 least 0.15) should be well-captured; reanalysis models use observed ice concentration in their 247 surface boundary condition, so the ice is imprinted onto their near-surface fields and then 248 recreated in the forced ocean model. However, hindcasting ice area (the area integral of ice 249 concentration) and thickness, hence volume, is more challenging. Antarctic ice drift is 250 dominated by surface winds, and ERA-Interim is known to capture the appropriate wind 251 252 trends (Holland and Kwok 2012). ERA-Interim air temperatures (Bracegirdle and Marshall 2012) and our model ocean temperatures (see below) are also reasonable, implying little 253 limitation on the ice hindcast. However, ice concentration and (crucially) thickness are 254 255 strongly affected by snow cover (Powell et al. 2005) and ocean freshwater fluxes (Hellmer 2004; Zhang 2007), both of which are limited by the large uncertainty in reanalysis 256 precipitation fields (although ERA-Interim is among the best, according to Bromwich et al. 257 (2011)). Also, any convergence-driven dynamical ice thickening will be determined by the 258 assumed rheology of the ice, of which model treatments are uncertain (Feltham 2008; 259 Tsamados et al. 2013). We therefore expect modelled ice thickness trends to be affected by 260 poorly-constrained details of the forcing and models. As a result, we perform a qualitative 261 assessment of our model results against existing observations, and consider broad patterns of 262 263 ice thickness change rather than quantitative predictions for specific regions, which are perhaps better-provided by the data-assimilating model of Massonnet et al. (2013). 264

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266 **3.1 Modelled Ocean Mean State**

Since the ocean state and trends can potentially have a significant effect on sea ice, we first assess the mean state of our modelled ocean over the period of interest, 1992-2010 (Figure 1).

The long-term mean barotropic streamfunction of the model (Figure 1a) reproduces the observed path of the Antarctic Circumpolar Current (Orsi et al. 1995; Sokolov and Rintoul 2009) and, crucially, also captures the shape and strength of the subpolar Weddell and Ross gyres (Wang and Meredith 2008). Thus, to the extent permitted by the sparse available data, we can have some confidence that the dynamic coupling between ocean and ice is accurate.

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The thermodynamic interaction is harder to verify, since there are very few relevant 275 observations of the ocean beneath Antarctic sea ice. Most of our knowledge of ice-ocean 276 277 interaction comes from summertime observations of the remnant Winter Water and shelf waters formed by winter sea-ice production. As summarised by Petty et al. (submitted), these 278 observations show that in the Weddell and Ross seas the surface mixed layer extends to the 279 280 seabed in winter, filling the shelf seas with cold and saline shelf waters, while in the Amundsen and Bellingshausen seas the winter mixing only produces a shallower layer of 281 Winter Water, beneath which warmer Circumpolar Deep Water is allowed to persist on the 282 shelf. The mean winter mixed-layer depth (Figure 1b) predicted by the KPP scheme (defined 283 as the shallowest depth for which the overlying bulk Richardson number equals 0.3) shows 284 that the model is able to reproduce these features, with complete destratification in the 285 Weddell and Ross seas and progressively shallower convection in the Amundsen and 286 Bellingshausen seas. This is also reflected in the long-term mean temperature and salinity at 287 288 the seabed, which shows warm and relatively fresh Circumpolar Deep Water in the Amundsen and Bellingshausen seas and cold and saline shelf waters in the Weddell and Ross 289 seas (Figures 1c and 1d). Further offshore, the winter mixed layer shallows over the sea-ice 290 zone due to a reduction in surface stress and buoyancy forcing, and then deepens offshore of 291 the ice edge. Thus, the vertical structure of the water column seems to compare well to the 292

limited observations that exist, and we infer that the thermodynamic ice-ocean interaction isreasonable as far as it can be tested.

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296 **3.2 Modelled Ice Mean State**

We next compare the mean state of our modelled Antarctic sea ice to observations over the 297 period of interest, 1992-2010. A comparison of mean ice concentration by season (Figure 2) 298 shows that the modelled ice concentration in Austral autumn and winter are very good, which 299 is critical because these seasons have the largest observed ice concentration trends (Turner et 300 301 al. 2009) and are best covered by ice motion data. Concentrations in spring and summer are not as good, with two persistent problems. Firstly, the model fails to capture a 'halo' of low 302 ice concentration near 0°E in spring (Lindsay et al. 2004), which leads to excessive summer 303 304 ice concentration in the eastern Weddell Sea. The halo is thought to be caused by upward 305 deformation of warm isopycnals near the Maud Rise seamount (de Steur et al. 2007), which is a challenging feature to capture accurately in a large-scale z-level ocean model. Attempts 306 307 were made to produce this feature using a variety ocean mixing schemes, but these resulted in open-ocean convection and a large polynya in the region (Timmermann and Beckmann 308 2004), strongly degrading the agreement with observations. Secondly, low ice concentrations 309 in the Ross Sea polynya are poorly represented in both spring and summer. Northward ice 310 export in this region is reasonable (see below), so this problem is due to excessive 311 312 importation of ice from the east.

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A similar comparison of effective ice thickness (Figure 3) shows reasonable results, although some ice concentration errors are also apparent in effective thickness. The model captures the general magnitude of ice thickness and correctly produces thicker ice in the Weddell, Bellingshausen, and Amundsen seas, though the spatial patterns within each region are

imperfect. The model under-represents the thickest ice in the north-west corner of the 318 Weddell Sea, though this problem is minimal in autumn, the season of greatest interest here. 319 Ice is too thick in the eastern Weddell Sea, in accordance with the aforementioned lack of 320 halo in this region, and the model over-predicts effective ice thickness in the Ross Sea 321 polynya in all seasons. A similar validation of effective snow thickness (Figure 4) is perhaps 322 worse, with the model failing to reproduce the correct thicknesses in summer and autumn, 323 and producing the wrong pattern in the Weddell Sea in spring. This is unsurprising given the 324 uncertainty surrounding reanalysis precipitation fields, but does place a limitation on our 325 326 results because snow flooding is an important component of Antarctic sea ice growth (Powell et al. 2005). The model produces a relatively good representation of effective snow thickness 327 in the Pacific sector in spring. 328

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330 Given the model's better performance in autumn and winter, and the larger ice trends and greater availability of data in those seasons, the rest of this study concentrates primarily on 331 those seasons. Figure 5 shows the mean ice velocities predicted by the model, which agree 332 with the observations rather well. The focussed northward ice export from the Ross Sea and 333 widespread export in the Weddell Sea are reproduced well, as is the westward coastal current 334 around East Antarctica. Ice drift is a little too rapid near coastlines and the ice edge. This 335 may be a feature of the coarse sampling of the ice observations in these regions, but is more 336 337 likely to be inaccuracy in the modelled ice dynamics (Uotila et al. submitted). Near the coast this could be caused by problems with the coarse wind forcing or ice rheology. The over-338 zealous coastal current in the Pacific sector transports too much ice from the Bellingshausen 339 340 and Amundsen seas into the Ross Sea, explaining the excessive ice concentration in the latter.

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The modelled seasonal cycle in total Antarctic ice area (the area integral of ice concentration) 342 compares extremely well with observations (Figure 6a), which is an important result because 343 ice area is much harder to reproduce in a model than ice extent. The mean cycle of total 344 Antarctic ice volume (Figure 6b) is also in excellent agreement with the data that exist. The 345 modelled Antarctic-wide average ice thickness (total ice volume divided by total ice area; 346 Figure 6b) is remarkably constant throughout the year, varying by less than 20%. This 347 implies that autumn/winter ice thickening is offset by the growth of large areas of thin ice, 348 and spring/summer ice thinning is offset by the melting of large areas of thin ice. The 349 350 observations suggest the possibility that thicker ice in summer is missed by the model, but this is uncertain because the observations are derived with different assumed values for snow 351 density in each season. If a uniform snow density were used for all seasons, the derived ice 352 353 thickness would be larger in spring and smaller in summer, in closer agreement with the model. 354

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356 **3.3 Modelled Ice Trends**

Figure 6 also provides an overview of modelled and observed trends in Antarctic sea ice. 357 Monthly anomalies of ice area from the mean seasonal cycle for the respective datasets are 358 remarkably consistent between model and observations (Figure 6c), with a few exceptions, 359 leading to a good prediction of the overall magnitude of the area trend. Given the difficulty 360 361 inherent in hindcasting ice area, this is an encouraging result that leads to some confidence in the modelled trends. Building on this confidence, Figure 6d shows a primary conclusion of 362 this model study, that overall Antarctic ice volume and Antarctic-wide average ice thickness 363 have both increased over 1992-2010. The overall volume increase of 29 km³/y is in good 364 agreement with the central Massonnet et al. (2013) estimate of 36 km³/y for 1980–2008 365 using data assimilation. The ice volume anomaly timeseries largely follows that of ice area 366

(Figure 6c), but there are several occasions where anomalies in average ice thickness 367 contribute significantly to ice volume, such as in the prolonged negative anomaly in both 368 variables between 2002 and 2004. As fractions of their mean annual values, the increases in 369 Antarctic average ice thickness (1.5 mm/y / 0.7 m \sim 0.2 %/y) and total area (20×10³ km²/y / 370 10^7 km² ~ 0.2 %/y) contribute equally to the trend in ice volume (30 km³/y / 7×10³ km³ ~ 0.4 371 %/y). The Antarctic average ice thickness trend produces a feasible increase of 2.6 cm over 372 the period considered. It is noteworthy that the simulation of Zhang (in press) produces a 373 similar fractional trend in ice volume (0.46 %/y) despite having ice that is approximately 374 375 double the observed thickness; this suggests that the Zhang (in press) thickness and volume trends are approximately twice the real value (since the extent trend is accurate in that study). 376

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378 These overall timeseries hide a strong pattern of regional variation in the trends, much of which compensate, so that the overall Antarctic-mean trends are the residual of much larger 379 regional changes. Figure 7 compares, by season, the maps of linear trend in modelled and 380 observed ice concentration. The general agreement is exceptionally good, with the model 381 clearly reproducing the wave-like pattern of ice concentration trends during this period: 382 decreasing ice cover in Bellingshausen, Weddell, and Mawson seas, and increasing ice cover 383 in Ross, Amundsen, and Cosmonaut seas (Holland and Kwok 2012). The model trends are 384 least reliable in summer, which is unsurprising given the above validation of mean ice 385 386 concentration in this season. It is interesting to note that the modelled concentration trends seem to be shifted eastwards relative to the observed trends. We cannot be sure why this is, 387 but speculate that the reanalysis winds place the climatological lows in the circumpolar 388 pressure trough (and thus their trends) too far east as a result of poorly representing the 389 deepening of low pressure systems as they navigate Antarctic topography. 390

We again restrict our attention to autumn and winter, and investigate the agreement of trends 392 in ice drift between model and observations (Figure 8). Since ice thickness is strongly 393 affected by convergence and divergence, it is essential to have confidence in our modelled 394 ice-drift trends if we are to believe our modelled thickness trends. As shown in Figure 8, the 395 dynamical trends in autumn are in good agreement with observations, particularly 396 considering how challenging it is to correctly model ice velocities, let alone their linear trend. 397 This agreement is largely the result of accurate surface wind trends in ERA-Interim (Holland 398 and Kwok 2012). In autumn the model correctly produces the observed decadal increase in 399 400 northward ice export in the Ross, Amundsen, and Cosmonaut seas, and the observed decrease in northward export in the Weddell and Mawson seas (Holland and Kwok 2012). Wind and 401 ice-dynamical trends in winter do not fit the observations quite as well, but the broad features 402 403 of a southward trend in the Bellingshausen Sea and northward trend around 0°E are found in both model and observations. 404

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406 These observational assessments of modelled trends in ice concentration and velocity allow us to critically consider the pattern of trends in effective ice thickness (Figure 9) that cause 407 the overall increase in Antarctic sea ice volume. It is immediately apparent that the regional 408 trends in effective ice thickness are at least an order of magnitude larger than the Antarctic-409 mean trend (Figure 6), which is their residual. The largest effective thickness trends (up to 5 410 411 cm/y) are found in the Amundsen Sea in winter. This further demonstrates that, while overall Antarctic ice trends may be subtle, the local changes can be of a considerable magnitude. 412 Unsurprisingly, we find that around the ice edge the spatial distribution of effective ice 413 thickness trends (Figure 9) mimics the trends in ice concentration (Figure 8), although there 414 are differences in the relative magnitude of these trends. More importantly, the model also 415 produces effective ice thickness trends in the internal ice pack near the coast, which are not 416

417 apparent in the concentration trends because the ice is close to full cover throughout the 418 periods considered. These 'internal' thickness trends have the largest regional magnitudes, 419 and are an important finding of this study. There are three main regions of internal ice 420 thickness increase: the northwest Weddell Sea in autumn, southern Weddell sea in autumn 421 and winter, and the Amundsen and Bellingshausen seas in winter. Similar trends appear in 422 the model results of Massonnet et al. (2013) and Zhang (in press), though their seasonal 423 structure and physical origin have not been fully examined.

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425 It is important to note that the maps of trend in effective ice thickness (volume ice per area ocean) are nearly identical to maps of trend in average ice thickness (volume ice per area ice). 426 Away from the ice edge the concentration remains near full cover throughout, so the effective 427 428 and average thickness are practically the same. Near the ice edge the average ice thickness is of order 10 cm, so the observed changes in ice concentration alone, of order 1%/y, would 429 give a change in effective ice thickness of order 1 mm/y. This is negligible compared to the 430 431 modelled effective ice thickness changes of order 1 cm/y, which are therefore demonstrated to be the result of large changes in average ice thickness. In other words, the trends in 432 effective ice thickness (volume ice per area ocean) near the ice edge in Figure 9 are 433 negligibly affected by the trends in ice concentration (area ice per area ocean) in Figure 8; 434 they are instead almost entirely trends in average ice thickness (volume ice per area ice). To 435 436 investigate these trends further we now consider a diagnostic decomposition of the icethickness equation. 437

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439 **3.4 Analysis of ice trends**

An overview of the processes governing the evolution of effective ice thickness can be obtained by separating the total tendency of effective thickness into dynamic and thermodynamic parts. Effective thickness is governed by a simple conservation equation

$$\frac{\partial h}{\partial t} = -\nabla . \left(\boldsymbol{u} h \right) + f$$

where h is effective thickness and u is velocity. The first term on the right-hand side is the 443 thickness change caused by ice-flux divergence, as determined by the momentum balance, 444 while f is the change in thickness due to thermodynamic processes. We record the values of 445 each of these terms separately, and the total tendency, in the ice code. Examination of the 446 mean fields of these tendency terms is highly instructive, as shown in an observational 447 assessment by Holland and Kwok (2012), but here our purpose is to assess trends in ice 448 thickness, for which we assess trends in the tendency terms. This analysis is performed for 449 autumn and winter only, the seasons for which we have greatest confidence in the model 450 451 results.

452

As with all such calculations, maps of interannual trend in the tendency terms are generated 453 by constructing seasonal means of the terms at each grid point and then calculating the 454 interannual trend in the values for each season. The tendency terms represent the rate of 455 456 change of effective ice thickness during a particular season, so our calculated trends represent the change in that rate over the decadal time period considered. For this reason, the trends in 457 effective ice thickness (e.g. Figure 10a) do not exactly match trends in effective ice thickness 458 459 tendency (e.g. Figure 10b). The former is the trend in mean autumn effective ice thickness, while the latter is the trend in the mean change in effective ice thickness over autumn. For 460 example, some of the trends in autumn ice thickness are caused by thicker ice being present 461 462 at the end of summer, and this would cause the two quantities to disagree. However, trends in the autumn effective thickness tendency (Figure 10b) do explain many of the features in 463

the autumn effective thickness trend map (Figure 10a). The only significant regions of disagreement are the areas of ice thickness increase in the southern Amundsen and Ross seas and northwest Weddell Sea, which are therefore revealed to be the result of summertime trends. The model performance is imperfect in summer, so these features should be treated with caution.

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The majority of the trends in effective ice thickness (Figure 10a) are reflected in the trends in 470 effective thickness tendency (Figure 10b), which we can decompose exactly into dynamic 471 472 (Figure 10c) and thermodynamic (Figure 10d) parts. This decomposition reveals that the trends in the Pacific sector are mostly explained by changes in ice dynamics (compare 473 Figures 10b and 10c). The autumn thickness trend in the southern Weddell Sea is also caused 474 475 by dynamics, but the thinning in the northern Weddell Sea, and most of the changes around East Antarctica, are due to thermodynamic changes. Changes in wind stress (Figure 10a) 476 succinctly explain all of these changes. In the Amundsen and Ross seas, increased northward 477 478 ice transport in autumn causes thinning in the south and thickening in the north. In the Bellingshausen Sea, a southward trend in wind stress causes the exact opposite, a loss of ice 479 from the ice edge and a strong thickening near the coast. In the Weddell Sea a decrease in 480 northward ice export away from the coast causes strong thickening. The thermodynamic ice 481 loss to the north could be a result of the decreased export of cold and dry air from Antarctica, 482 or perhaps a southward shift of the warmer waters of the ACC, either of which could be 483 caused by the wind trends. The remaining trends all follow the same pattern of increased 484 (decreased) northward wind stress causing ice thickness increase (decrease) near the ice edge, 485 through a varying combination of changes in air-ice drag and cold- or warm-air advection. 486 These results are in complete agreement with the analysis of Holland and Kwok (2012), who 487 used observations to perform an autumn decomposition of the conservation equation for ice 488

concentration. Wind-driven ice convergence and a resultant thickening in the Pacific sector and southern Weddell Sea were also obtained by Zhang (in press). Finally, we note that the decomposition suggests an increased ice divergence and thermodynamic ice growth in the Ross Sea coastal polynya (Figures 10c and 10d), and a decrease in divergence and growth in the Ronne polynya, Weddell Sea, both in agreement with observed trends (Drucker et al. 2011).

The results in winter (Figure 11) illustrate the difference between trends in effective ice 496 497 thickness and effective ice thickness tendency. In this season few of the large ice thickness trends (Figure 11a) are observed in the tendency terms (Figure 11b), implying that the 498 thickness trends are the result of changes occurring in previous seasons. For example, the ice 499 500 thinning trend in the northern Weddell Sea (Figure 11a) is revealed as being a lasting effect of previous seasons; the trend in winter tendency (figure 11b) is towards thickening. On 501 average, there is thinner ice in the northern Weddell Sea during winter, but this ice is 502 503 thickening more during winter. The ice is thickening less during autumn, and the ice remains thinner during winter as a result. The increased thickening during winter is revealed as being 504 505 dynamical in origin (figure 11c), because the wind trend in this region is towards increased northward flow (figure 11a). 506

507

Some effective thickness trends that are very clearly caused by wintertime changes are in the Bellingshausen and Amundsen seas, where strong wind trends towards the south lead to a significant winter thickening of the ice near the coast that is entirely dynamic in origin (Figure 11). It is virtually certain that these thickening trends have occurred in reality, since they are the logical extension of known trends in ice concentration, winds, and ice drift in this region (Holland and Kwok 2012; Turner et al. 2009). The model is clearly responding

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sensibly to the wind stress it receives from ERA-Interim (Figure 11a). However the 514 magnitude and pattern of this thickening must be regarded as merely indicative, for two 515 reasons. Firstly, the ice model cannot be expected to convert wind stress changes into ice 516 thickness changes with a high level of quantitative skill, because this process is heavily 517 dependent upon the poorly constrained rheological properties of the ice (Feltham 2008; 518 Tsamados et al. 2013). Secondly, the detailed pattern of the southward trend in modelled ice 519 motion in this region in winter is imperfect (Figure 8); the observed ice drift trend is towards 520 the Antarctic Peninsula, while the reanalysis wind stress trend (Figure 11a) drives the ice 521 522 towards the coast in the eastern Bellingshausen Sea and the Amundsen Sea. However, ice drift trends in autumn are well-represented (Figure 8), and these do drive ice westwards in the 523 In summary, the observations strongly support a significant coastal ice 524 observations. 525 thickening in this region, but the model may place it too far east, and with an uncertain magnitude. Massonnet et al. (2013) also model a narrow zone of coastal thickening in this 526 region; Zhang (in press) does not. 527

528

529 4. Discussion

The model results presented here reproduce observations of mean ice concentration, drift, and 530 thickness, and trends in ice concentration and drift. The simulated ice thickness trends also 531 agree with those of Massonnet et al. (2013), which can be regarded as a 'best estimate' due to 532 533 their use of data assimilation. This gives us confidence that the physical processes in the model reflect those operating in reality, offering insight into the processes causing trends in 534 Antarctic sea ice. Holland and Kwok (2012) showed that autumn ice concentration trends are 535 536 dominated by dynamics in the Pacific sector of the Southern Ocean and thermodynamics elsewhere; this modelling study shows that the same pattern holds for ice thickness, and 537 hence ice volume, in autumn and winter. 538

This finding has significant consequences. Ice dynamical changes can occur either because 540 the driving stresses have changed, or because the ice is responding differently to a constant 541 stress. The latter can occur if the ice thins, since weaker ice responds more readily to an 542 applied stress, and this is the case in the Arctic, where the ice is accelerating in excess of 543 trends in wind forcing (Kwok et al. 2013). In the Antarctic the trends in ice motion and wind 544 agree closely (Holland and Kwok 2012) and the thickness changes modelled here are much 545 smaller. Thus, the ice-dynamical changes can only ultimately be caused by changes air-ice 546 547 drag and/or ocean-ice drag, which both ultimately result from changes in the winds since surface ocean currents are predominantly wind-driven. The dynamic origin of the modelled 548 changes in the Pacific sector in autumn and winter therefore implies little or no contribution 549 550 from changes due to precipitation (Liu and Curry 2010), feedbacks (Stammerjohn et al. 2012; Zhang 2007), or atmosphere or ocean warming (Jacobs and Comiso 1997; Lefebvre and 551 Goosse 2005; Liu et al. 2004). This certainly does not rule out a contribution from these 552 mechanisms in summer and spring, or around East Antarctica. A detailed analysis of the 553 trends in ice thermodynamics, in a model capable of accurately representing the warmer 554 seasons, is clearly required to advance this question. 555

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The results also suggest that it is unlikely that increased ice-sheet melting is implicated in the Antarctic sea ice increase, as proposed by Bintanja et al. (2013). The vast majority of increased freshwater discharge from the Antarctic Ice Sheet has entered the Amundsen Sea (Shepherd et al. 2012) and followed the coastal current westward into the Ross Sea, where it has caused a significant freshening (Jacobs and Giulivi 2010). If ice-shelf meltwater were to contribute to the sea-ice trends, the largest effect would thus be expected to occur in the increasing ice volume in the western Pacific. Our results, and the observational analysis of Holland and Kwok (2012), show quite clearly that the trends in that region are predominantly dynamic in origin in autumn and winter. In addition, the model presented here has no overall trend in ice-sheet meltwater input (the prescribed iceberg discharge is steady, and total iceshelf melting contains no significant trend), yet is able to reproduce most features of the observed Antarctic ice concentration increase. Thus our results are in agreement with the study of Swart and Fyfe (2013), who found that the Antarctic sea ice trends were not affected by trends in Antarctic Ice Sheet freshwater flux.

571

572 **5. Conclusions**

There are no observations of decadal trends in Antarctic sea ice thickness and volume, so we hindcast them for the period 1992—2010 using a numerical ice—ocean model that is extensively validated against observations. The model accurately simulates mean fields of ice concentration, drift, and thickness in autumn and winter, and reproduces observed trends in ice concentration and drift. This validation allows us to hold some confidence in the corresponding modelled trends in ice thickness.

579

Unsurprisingly, the model shows that the observed ice-concentration trends near the ice edge 580 have corresponding trends in ice thickness, with areas of increasing thickness associated with 581 increasing concentration. Model diagnostics show that these thickness trends are driven 582 583 dynamically in the Pacific sector and thermodynamically elsewhere, in agreement with an observational decomposition of ice concentration trends (Holland and Kwok 2012). The 584 model also reveals that the observed southward trends in ice drift in the Bellingshausen and 585 586 Weddell seas have caused ice to thicken near the coast, a trend that does not appear in ice concentration measurements because the ice remains at full cover throughout. The Weddell 587 Sea thickening occurs in response to decreased export early in the year, while the 588

Bellingshausen Sea thickening occurs in winter due to a strong trend towards southward ice flow. These results are the logical extension of known trends in ice concentration, winds, and ice drift. The dynamic origin of the autumn and winter trends in the Pacific sector imply that they must be forced by changes in the winds, rather than other atmospheric or oceanic forcings or feedbacks.

594

Spatial patterns of increasing and decreasing trends in ice concentration and thickness largely 595 compensate, so that neither variable has a large Antarctic trend overall. Thickening in the 596 597 interior of the ice pack enhances the overall thickness trend relative to the concentration trend. As fractions of their mean annual values, the modelled increases in Antarctic-wide ice 598 thickness (1.5 mm/y ~ 0.2 %/y) and area (20×10^3 km²/y ~ 0.2 %/y) contribute equally to the 599 overall trend in ice volume (30 km³/y \sim 0.4 %/y). This small gain contrasts markedly with 600 the observed Arctic sea ice volume loss of 500-1000 km³/y (~3-6 %/y) (Kwok and 601 Rothrock 2009; Laxon et al. 2013). In terms of Southern Ocean freshwater forcing, the small 602 increase in sea ice freshwater extraction is outweighed by the $\sim 70 \text{ km}^3/\text{y}$ increase in 603 freshwater input from the Antarctic Ice Sheet (Shepherd et al. 2012). 604

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784 Figure Captions

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Figure 1: Modelled mean 1992-2010 ocean fields. a) Barotropic stream function (contours
every 10 Sv, magenta contour 0 Sv), b) Winter (JJA) mean mixed-layer depth from KPP
calculation (contours every 25 m, magenta contour 100 m), c) Potential temperature at seabed
(contours every 0.2 °C, magenta contour 0 °C), d) Salinity at seabed (contours every 0.025,
magenta contour 34.65).

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Figure 2: Modelled and observed 1992-2010 mean ice concentration by season. Observed ice
concentration is calculated using the Bootstrap algorithm (Comiso 2000). WS: Weddell Sea,
CS: Cosmonaut Sea, MS: Mawson Sea, RS: Ross Sea, AS: Amundsen Sea, BS:
Bellingshausen Sea.

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Figure 3: Modelled mean 1992-2010 effective ice thickness and observed mean 2003-2008 effective ice thickness by season. Effective ice thickness is defined as volume of ice per unit area of ocean, neglecting the ice-borne snow layer. Observed effective ice thickness is derived from ICESat measurements (Kurtz and Markus 2012). Areas with respective ice concentration below 0.5 are masked in both datasets.

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Figure 4: Modelled mean 1992-2010 effective snow thickness and observed mean 2003-2008
effective snow thickness by season. Effective snow thickness is defined as volume of iceborne snow per unit area of ocean. Observed effective snow thickness is derived from
ICESat measurements (Kurtz and Markus 2012). Areas with respective ice concentration
below 0.5 are masked in both datasets.

Figure 5: Modelled and observed 1992-2010 mean ice concentration and ice drift for autumn
and winter seasons. Observed ice concentration is calculated using the Bootstrap algorithm
(Comiso 2000) and ice velocities are from passive microwave feature-tracking (Holland and
Kwok 2012). Model velocities are shown every tenth grid point.

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Figure 6: 1992-2010 temporal variability of total Antarctic sea ice variables from model and 814 observation. a) mean seasonal cycle in total ice area (the area integral of ice concentration) 815 816 for model and observations (Comiso 2000). b) mean seasonal cycle in total ice volume and mean ice thickness (total ice volume divided by total ice area) for model and observations 817 (dots represent individual ICESat campaigns, shaded areas represent interannual mean ± 818 819 standard deviation for each season; Kurtz and Markus (2012)). c) monthly anomalies in 820 modelled and observed total ice area from respective climatologies in panel a. d) monthly anomalies in modelled total ice volume and mean ice thickness from respective climatologies 821 in panel b. All trends shown are significant at the 99% level. 822

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Figure 7: Modelled and observed 1992-2010 linear trends in ice concentration by season.
Observed ice concentration is calculated using the Bootstrap algorithm (Comiso 2000).

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Figure 8: Modelled and observed 1992-2010 linear trends in ice concentration and drift for autumn and winter seasons. Observed ice concentration is calculated using the Bootstrap algorithm (Comiso 2000) and ice velocities are from passive microwave feature-tracking (Holland and Kwok 2012). Model velocities are shown every tenth grid point.

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Figure 9: Modelled 1992-2010 linear trends in effective ice thickness and drift for autumn
and winter seasons. Effective ice thickness is defined as volume of ice per unit area of ocean,
neglecting the ice-borne snow layer. Model velocities are shown every tenth grid point. The
largest trends, up to 5 cm/y, are in the Amundsen Sea in winter.

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Figure 10: Modelled autumn (AMJ) 1992-2010 linear trends in effective ice thickness and
related quantities. a) trends in modelled effective ice thickness and ERA-Interim wind stress
(shown every tenth grid point); b) trends in evolution term in ice-thickness equation; c) trends
in dynamic part of ice-thickness evolution; d) trends in thermodynamic part of ice-thickness
evolution. The colourbar for panels c and d is the same as for panel b.

Figure 11: As Figure 10 but for winter (JAS).



Figure 1: Modelled mean 1992-2010 ocean fields. a) Barotropic stream function (contours
every 10 Sv, magenta contour 0 Sv), b) Winter (JJA) mean mixed-layer depth from KPP
calculation (contours every 25 m, magenta contour 100 m), c) Potential temperature at seabed
(contours every 0.2 °C, magenta contour 0 °C), d) Salinity at seabed (contours every 0.025,
magenta contour 34.65).



Figure 2: Modelled and observed 1992-2010 mean ice concentration by season. Observed ice
concentration is calculated using the Bootstrap algorithm (Comiso 2000). WS: Weddell Sea,
CS: Cosmonaut Sea, MS: Mawson Sea, RS: Ross Sea, AS: Amundsen Sea, BS:
Bellingshausen Sea.



Figure 3: Modelled mean 1992-2010 effective ice thickness and observed mean 2003-2008 effective ice thickness by season. Effective ice thickness is defined as volume of ice per unit area of ocean, neglecting the ice-borne snow layer. Observed effective ice thickness is derived from ICESat measurements (Kurtz and Markus 2012). Areas with respective ice concentration below 0.5 are masked in both datasets.



Figure 4: Modelled mean 1992-2010 effective snow thickness and observed mean 2003-2008
effective snow thickness by season. Effective snow thickness is defined as volume of iceborne snow per unit area of ocean. Observed effective snow thickness is derived from
ICESat measurements (Kurtz and Markus 2012). Areas with respective ice concentration
below 0.5 are masked in both datasets.



Figure 5: Modelled and observed 1992-2010 mean ice concentration and ice drift for autumn
and winter seasons. Observed ice concentration is calculated using the Bootstrap algorithm
(Comiso 2000) and ice velocities are from passive microwave feature-tracking (Holland and
Kwok 2012). Model velocities are shown every tenth grid point.



Figure 6: 1992-2010 temporal variability of total Antarctic sea ice variables from model and 872 observation. a) mean seasonal cycle in total ice area (the area integral of ice concentration) 873 for model and observations (Comiso 2000). b) mean seasonal cycle in total ice volume and 874 mean ice thickness (total ice volume divided by total ice area) for model and observations 875 (dots represent individual ICESat campaigns, shaded areas represent interannual mean ± 876 standard deviation for each season; Kurtz and Markus (2012)). c) monthly anomalies in 877 modelled and observed total ice area from respective climatologies in panel a. d) monthly 878 anomalies in modelled total ice volume and mean ice thickness from respective climatologies 879 in panel b. All trends shown are significant at the 99% level. 880



Figure 7: Modelled and observed 1992-2010 linear trends in ice concentration by season.

883 Observed ice concentration is calculated using the Bootstrap algorithm (Comiso 2000).



Figure 8: Modelled and observed 1992-2010 linear trends in ice concentration and drift for autumn and winter seasons. Observed ice concentration is calculated using the Bootstrap algorithm (Comiso 2000) and ice velocities are from passive microwave feature-tracking (Holland and Kwok 2012). Model velocities are shown every tenth grid point.



Figure 9: Modelled 1992-2010 linear trends in effective ice thickness and drift for autumn
and winter seasons. Effective ice thickness is defined as volume of ice per unit area of ocean,
neglecting the ice-borne snow layer. Model velocities are shown every tenth grid point. The
largest trends, up to 5 cm/y, are in the Amundsen Sea in winter.



Figure 10: Modelled autumn (AMJ) 1992-2010 linear trends in effective ice thickness and related quantities. a) trends in modelled effective ice thickness and ERA-Interim wind stress (shown every tenth grid point); b) trends in evolution term in ice-thickness equation; c) trends in dynamic part of ice-thickness evolution; d) trends in thermodynamic part of ice-thickness evolution. The colourbar for panels c and d is the same as for panel b.



901 Figure 11: As Figure 10 but for winter (JAS).